

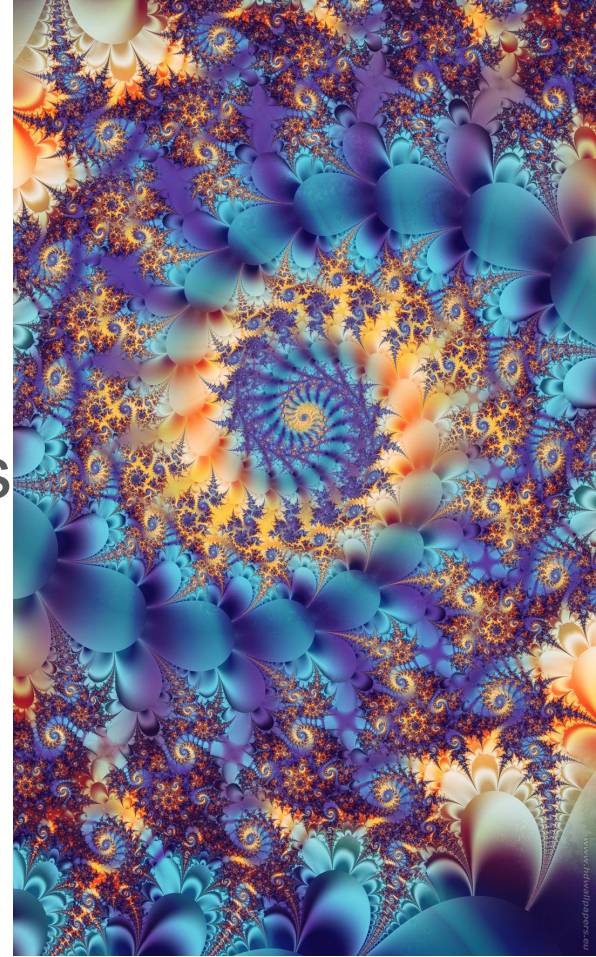
Computational Chaos by Edward N. Lorenz

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Spring 2017

Outline

- Introduction
- A system with fixed-point attractors
- Instability of fixed points
- Bifurcations
- Onset of chaos
- Strange Attractors
- Conclusion



Introduction

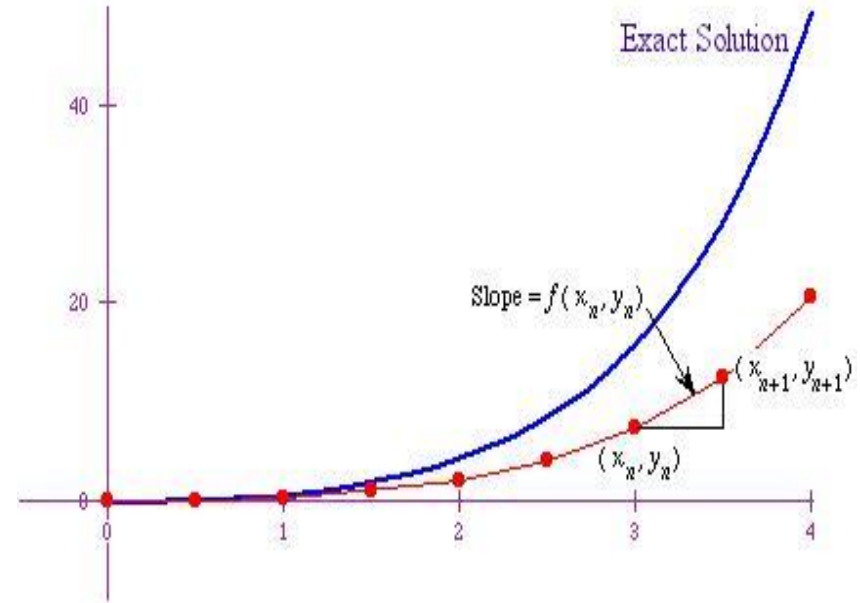
- Solutions to nonlinear differential equations often sought by numerical means.

$$d\mathbf{X}/dt = \mathbf{F}(\mathbf{X})$$

- Approximating using Euler scheme:

$$\mathbf{X}_{n+1} = \mathbf{X}_n + \tau \mathbf{F}(\mathbf{X}_n)$$

- Other techniques:
 - Runge-Kutta
 - 4th-Order Taylor Series



Introduction

- Chaotic behavior sometimes occurs when difference equations used as approximation to ordinary differential equation are solved numerically with a large time increment.
- Using one example we show when fixed points go unstable and when chaos first sets in

Introduction

- **Computational chaos:** A chaotic behavior that owes its existence to the use of a large time increment, τ .
- **Goal:**
 - Lower limit of values of τ for which fixed points are unstable?
 - Lower limit of values of τ for which chaos is present?

Introduction

- Computational chaos is widespread: Even for one of the simplest flows

$$dx/dt = x - x^2$$

- Almost all solutions approach either $-\infty$ or the stable fixed point $x = 1$
- If we apply Euler function:

$$x_{n+1} = (1 + \tau)x_n - \tau x_n^2$$

- $0 < x_0 < (1 + \tau)/\tau \Rightarrow$ fixed point $x = 1$ if $\tau < 2$
 \Rightarrow chaotic behavior if $2 < \tau < 3$

A system with fixed-point attractors

- Used a model of fluid convection and turned into a limiting form of it.

$$dX/dt = -\sigma X + \sigma Y,$$

$$dY/dt = -XZ + \rho X - Y, \quad \longrightarrow$$

$$dZ/dt = XY - \beta Z,$$

$$dx/dt = ax - xy,$$

$$dy/dt = -y + x^2$$

- As: $\sigma \rightarrow \infty$, $\beta = 1$, $a = \rho - 1$
- Replace X by Y
- Replace Y and Z by x and y

A system with fixed-point attractors

- Approximating the limiting form by Euler scheme:

$$x_{n+1} = (1 + a\tau)x_n - \tau x_n y_n$$

$$y_{n+1} = (1 - \tau)y_n + \tau x_n^2.$$

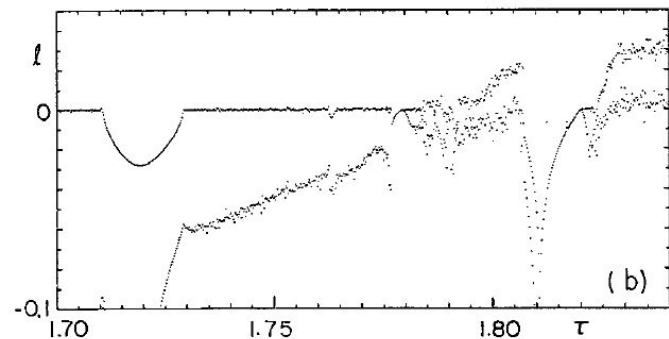
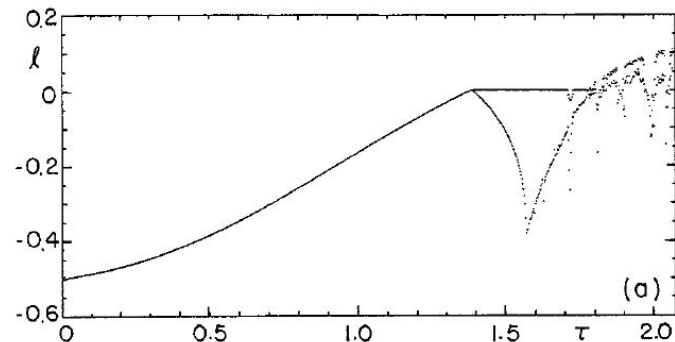
- After solving the differential equation we get Lyapunov exponents: l_1 and l_2 which are both equal to:

$$[\log(1 - \tau + 2a\tau^2)]/(2\tau).$$

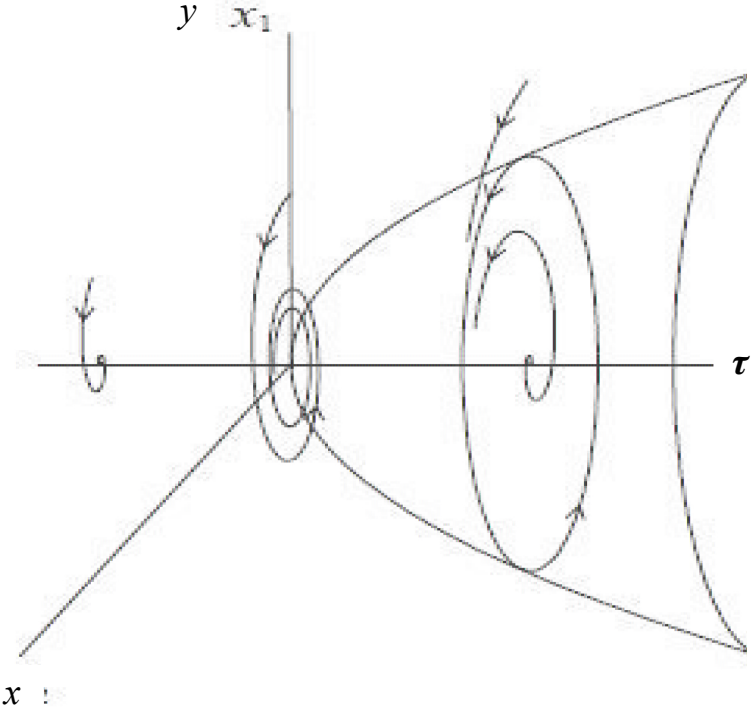
- If $\tau \rightarrow 0 \Rightarrow -1/2$ (stable fixed point)
- $\tau = 1/(2a) \Rightarrow 0$

Lyapunov Exponent VS Time Increment

- $0 \leq \tau < \tau_a, l < 0 \Rightarrow$ Stable fixed point
 - Single stable fixed point
- $\tau_a \leq \tau < \tau_b, l > 0 \Rightarrow$ Unstable fixed point
 - Hopf bifurcation
- $\tau_b \leq \tau < \tau_c, l > 0 \Rightarrow$ Unstable fixed point
 - Chaos is present
- $\tau_c \leq \tau$
 - Computational instability

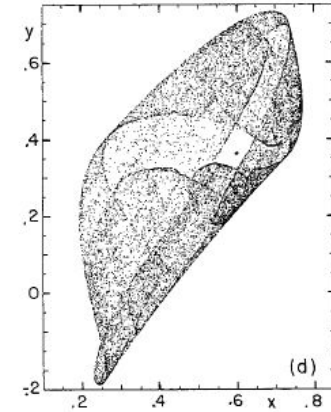
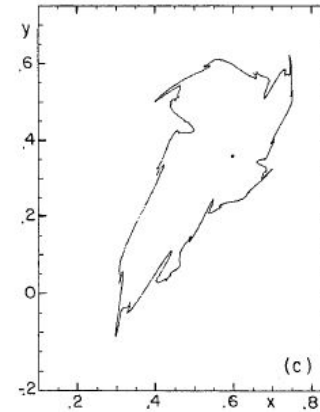
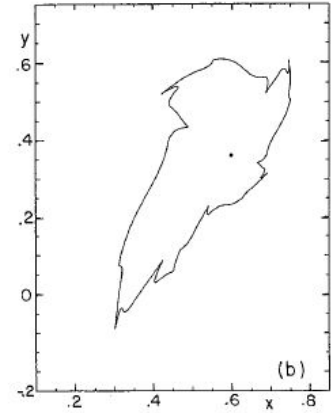
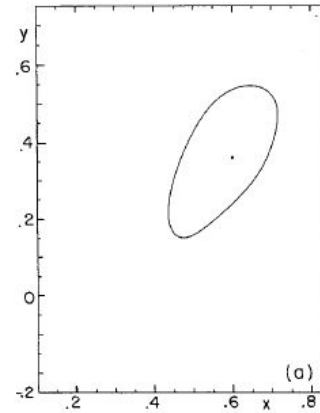


Hopf Bifurcation

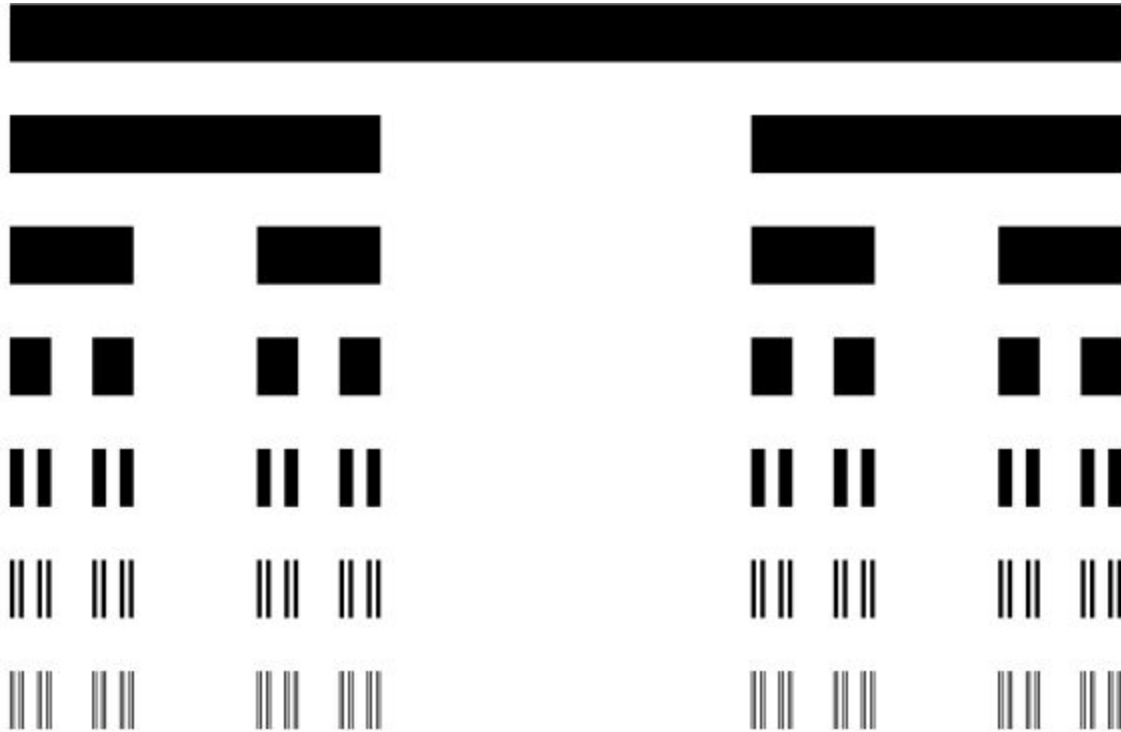


Attractors About A Fixed Point

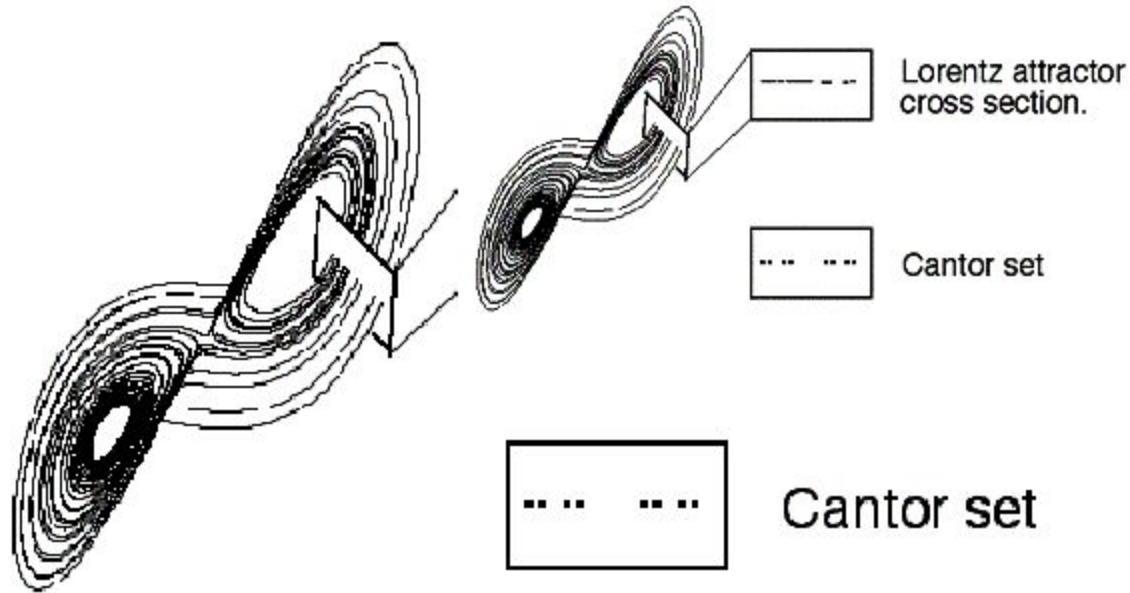
- Limit cycle at Q. Not chaotic
- Deformed limit cycle. Not chaotic
- Further deformed. Chaotic
 - Hopf Bifurcations at kinks



Cantor Set

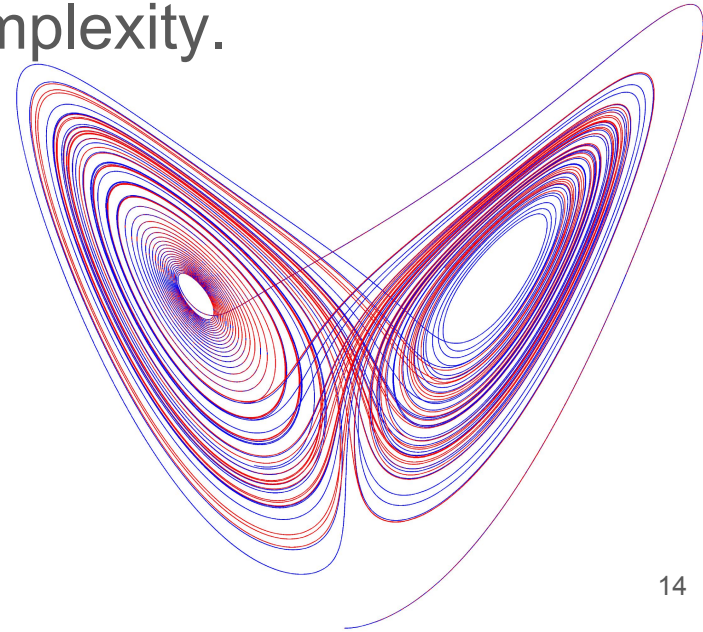


Lorenz Attractor Cross Section



Strange Attractor

- Closed set
- Can be visualized in phase space
- Fractal substructure with infinite complexity.
- Basin of attraction



Conclusion

- Interaction between parameters (a , b , τ) can lead to chaos
- Chaos starts with different types of bifurcation
- Not all approximation techniques are created equal
 - Runge-Kutta has a larger range of τ before chaos but can still go chaotic
- Strange attractors:
 - fractal “surfaces”
 - Resemble the Cantor set

References

- Lorenz, E., “Computational Chaos”, Physica D, 1988