Lecture 7

Order Out of Chaos

Lyapunov Exponents: Recall from Last Time

The Lyapunov exponent for maps is: $\lambda = \ln |f'(x^*)| = \begin{cases} < 0 \text{ if } |f'(x^*)| < 1 \\ > 0 \text{ if } |f'(x^*)| > 1 \end{cases}$

Lyapunov exponent: $\lambda = \ln |f'(x^*)|$

Rate of divergence over time:

 $f'(x) = \frac{dx_{t+1}}{dx_t}$

Lyapunov exponent: $\lambda = \ln |f'(x^*)|$

Rate of divergence over time:

 $f'(x) = \frac{dx_{t+1}}{dx_t}$

Assumption: Lyapunov exponent is the same everywhere in the basin of attraction.

Lyapunov Exponents: Maximum



This is really the averaged sum of the Lyapunov exponents, λ .

Lyapunov Exponents: Maximum

 λ : Short term behaviour λ_{\max} : Long term behaviour

Since maps depend on the previous

step, we can show that:

Mean rate of divergence:

$$\lambda_{\max} = \lim_{t \gg 1} \frac{1}{t} \ln \left| \frac{f_t(x)}{dx} \right| = \frac{1}{t_{\max}} \sum_{t=t_0}^{t_{\max}} \ln \left| \frac{dx_{t+1}}{dx_t} \right|$$

This is our approximation of λ_{\max}



We want to look at the behaviour near x^* . How can we know x_t is near x^* ?

We want to look at the behaviour near x^* . How can we know x_t is near x^* ?

We iterate the map long enough that x_t crosses the separatrix that defines the long term behaviour.

This is our approximation of λ_{\max}



Where t_{start} exceeds the transient period.



Code it up!

function ls = lyapunov(F, F_deriv, n_samples, param_range, transient_time, max_time)
% Takes a dynamical 1-parameter map and plots the Lyapunov exponents as a function of
% the parameter

```
% Create a vector of parameter values to evaluate
param_values = linspace(param_range(1), param_range(2), n_samples);
```

```
ls=[]; % This stores the Lyapunov exponents
```

```
for param=param values
    x t=rand(1); % Sample over many initial values
    lyp exps = [];
    for(t = 1:max time)
        % Evaluate the user defined map F
        x t=F(x t, param);
       % Wait until the transient period is over
       if(t > transient time)
           % Evaluate the derivative at the current point
           lyp exps = [lyp exps, F deriv(x t, param)];
       end
    end
```

```
% Calculate Lyapunov approximation for the vector of derivatives
ls = [ls, mean(log((abs(lyp_exps))))];
```

end

% Show the Lyapunov exponents and bifurcation plot for the logistic map F = @(x,r) r*x*(1-x) % Define the logistic map F_deriv =@(x,r) r-2*r*x % Define the derivative of the logistic map x0 = 0.5 % An initial value for the bifurcation plot n_samples = 250 % Number of points to plot param_range = [0,4] % Parameter range to plot, r for the logistic map transient_time = 500 % Make sure we are in the fixed point's basin max_time = 1000 % This minus the transient time is the number % of Lyapunov samples to average over

% Calculate the Lyapunov exponents
ls = lyapunov(F, F_deriv, n_samples, param_range, transient_time, max_time);

% Calculate the long term populations bs = bifurcation(F,x0,param_range(1),param_range(2),n_samples,max_time);

% Plot the results subplot(2,1,1) plot(linspace(param_range(1), param_range(2), n_samples), ls, 'b-'); subplot(2,1,2) plot(linspace(param_range(1), param_range(2), n_samples), bs, 'k.');



Top Panel: Plot of Lyapunov Exponents, Bottom Panel: Bifurcation Plot for r=0.5



Top Panel: Plot of Lyapunov Exponents, Bottom Panel: Bifurcation Plot for r=0.5

Lyapunov Exponents for 2D Maps

The Duffing or Holmes Map: $x_{t+1} = y_t$ $y_{t+1} = -bx_t + ay_t - y_t^3$ Chaotic at a = 2.75 and b = 0.2.

Lyapunov Exponents for 2D Maps

This time we have to deal with divergence in 2-dimensions. To do that we use the Jacobian.

See Lecture 5 for our discussion of the Jacobian.

```
current_1 = 0;
for param1=param1_range
    current_1 = current_1 + 1;
    % Initialize variables
    xy = [x0; y0]; xy_lengths = [1;0];
    for i=1:t_max
        J = F_Jacobian(xy, param1, param2);
        xy=F(xy, param1, param2);
```

end

end

% Calculate divergence rate in the direction defined by the Jacobian xy_lengths=J*xy_lengths; length=sqrt(sum(xy_lengths.^2)); % Distance formula max_lyapunovs(current_1) = log(length)/i; % Calculate the average end $F = @(xy,a, b) [xy(2); -b*xy(1)+a*xy(2)-xy(2)^3] % Duffling Map$ $F_Jacobian = @(xy,a, b) [0 1; -b a-3*(xy(2))^2] % Duffling Jacobian$

max_time = 500; % How long to run (= number of samples to average)
parameter1_range = 2:0.001:3; % Range over parameter 1: (a for Duffling Map)
parameter2 = 0.2; % Fix parameter 2 (b for Duffling Map)

```
% Initial values for x and y
x0 = 0.5
y0 = 0.5
```

% Calculate the maximum Lyapunov exponents max_lyapunovs = lyapunov2d(F, F_Jacobian, max_time, parameter1_range,... parameter2, x0, y0);

```
% Make a plot of the maximum exponents with a line at 0
plot(parameter1_range,max_lyapunovs, parameter1_range, 0, 'k.')
xlabel('Param1: a', 'FontSize', 24);
ylabel('\lambda_{max}', 'FontSize', 24);
```

Example: Holmes Map





Feigenbaum's Constant

$$\frac{R_1 - R_0}{R_2 - R_1} \approx 4.75$$

$$\frac{R_{n-1} - R_{n-2}}{R_n - R_{n-1}} \approx 4.669$$

Feigenbaum's Constant

True for ALL maps that approach chaos by bifurcation.

$$\frac{R_1 - R_0}{R_2 - R_1} \approx 4.75$$

$$\frac{R_{n-1} - R_{n-2}}{R_n - R_{n-1}} \approx 4.669$$

Phase Space Contraction

We can look at the evolution of a small volume (in 3D) as it contracts near an attractor The change in volume is measured by looking at the trace of the Jacobian, $\vec{\nabla} \cdot \dot{\vec{x}}$

A dynamical system is dissipative, if it's phase space volume contracts continuously,

 $\vec{\nabla} \cdot \dot{\vec{x}}(t) < 0, \forall t$



Dissapative or conservative?





Dissapative or conservative?

A dynamical system is conserving, if it's phase space volume is constant, $\vec{\nabla} \cdot \dot{\vec{x}}(t) = 0, \forall t$

Adaptive Systems

- Adaptive systems are neither fully
- dissipative nor fully conserving.
- In terms of energy they have periods
- of taking up energy and periods of expending it.
- Technically, any system where $\vec{\nabla} \cdot \dot{\vec{x}}(t)$ can change sign is adaptive.

Outline

- Feigenbaum Constant (Order -> Chaos -> Deeper Order)
- Kuhns Revolutions in Science
 - Microscope
 - Telescope
 - Computer (George Luger)
 - May, Lorenz, Crutchfield, Mitchell and many others...
- Biological Complexity from Simple Rules. Hearts (fractal structure)
- Exploring Dynamical Systems has demonstrated that surprising complexity can arise from simple rules.
- This allows us to understand the diversity generated by things like capitalism, biological evolution, computer architectures, and computer algorithms.